without correction for beam inhomogeneity; that is, to the same level of approximation as in Figure 4.

Three observations are worth making. (1) In the quantitative interpretation of data, beam inhomogeneity needs to be taken into account, especially at low dosages. (2) The empirical eq 17 is not an exact relationship, because neither of the curves in Figure 5A is a horizontal straight line. According to the dashed curve, the parameter A in (17) varies by a factor of ~ 2 while $\exp(-v_{act}/v_{abs})$ varies by a factor of ~3000. (3) All relationships shown in Figure 5 increase up to $v_{act} \sim 5$ and appear to become constant when $v_{act} > 5$.¹⁸ This feature is shown most clearly by the solid curve in Figure 5B, which is based on the exact harmonic oscillator, eq 20. The maximum value of f/x_{act} is about 5, which means that the amount of laser-induced reaction is greater than the fraction of activated molecules predicted by the simple harmonic oscillator model.

References and Notes

- (1) Work supported by the National Science Foundation.
- (a) M. W. Berry, Annu. Rev. Phys. Chem., 26, 263–268 (1975); (b) J. T. Knudtson and E. M. Eyring, *ibid.*, 25, 255 (1974).
 (3) W. M. Shaub and S. H. Bauer, Int. J. Chem. Kinet., 7₁ 509 (1975).
- (4) N. V. Karlov, Appl. Opt., 13, 301 (1974).

- (5) I. Glatt and A. Yogev, J. Am. Chem. Soc., 98, 7087 (1976).
- (6) D. F. Dever and E. Grunwald, J. Am. Chem. Soc., 98, 5055 (1976).
- E. Grunwald and K. J. Olszyna, Laser Focus, 12(6), 41 (1976). (8) R. G. Barnes, R. A. Cox, and R. F. Simmons, J. Chem. Soc. B, 1176
- (1971).
- (9) E. K. Plyler and W. S. Benedict, J. Res. Natl. Bur. Stand., 47, 212 (1951).
- (10) Synonymous with dose/flash are the terms "energy flux" or "energy fluence". We prefer to retain the traditional term "dose", which has been used for many years in such fields as X-radiology and radiation chemis-
- (11) (a) E. D. West and D. A. Jennings, Rev. Sci. Instrum., 41, 565 (1970); (b) R. W. Zimmerer, "Theory and Practice of Thermoelectric Laser Power and
- K. W. Zimmerer, "Indify and tradine of minimediate Data (2014) and the state of the sta
- (14) K. J. Olszyna, E. Grunwald, P. M. Keehn, and S. P. Anderson, Tetrahedron Lett., 1609 (1977)
- (15) R. N. Schwartz, Z. I. Slawsky, and K. F. Herzfeld, J. Chem. Phys., 20, 1591 (1952). (16) T. D. Rossing and S. Legvold, *J. Chem. Phys.*, **23**, 1118 (1955).
- (17) J. M. Preses, R. E. Weston, Jr., and G. W. Flynn, Chem. Phys. Lett., 46, 69 (1977).
- (18) Because f/x_{act} approaches unity in the limit as v_{act} becomes large, the crosses in Figure 5 represent the function ln (1 f)/ln (1 x_{act}), which does not approach any mathematically required limit as vact becomes large. beside the upper sector of the sector of th of f seems logical for expressing amount of reaction in an inquiry which probes the chemical kinetics.

Megawatt Infrared Laser Chemistry of CCl₂F₂ Derived from the Excitation of Two Distinct Vibrational Modes¹

Gregory A. Hill, Ernest Grunwald,* and Philip Keehn

Contribution from the Department of Chemistry, Brandeis University, Waltham, Massachusetts 02154. Received February 25, 1977

Abstract: CCl₂F₂ was pulse irradiated at 921 cm⁻¹, where a CCl₂ stretching mode is excited, and at 1088 cm⁻¹, where a CF₂ stretching mode is excited. Laser-induced reactions at the two frequencies are practically identical, both qualitatively and quantitatively. Major reaction products are CIF_2CCCIF_2 and $CCIF_3$ and account for ~83% of decomposition. Minor products formed initially are CIF₂CCCl₂F and CCl₃F; there was no evidence for C₂F₄. A reaction mechanism involving primary C-Cl bond scission is indicated. Conversion of CCl_2F_2 per flash (f) was as high as 13%; radiant energy absorbed (E_{abs}) was as high as 107 kJ/mol. The plot of ln f vs. E_{abs}^{-1} is linear with a slope of -419 kJ/mol, which is in good agreement with $-E_{act}$ for thermal decomposition of CCl₂F₂. When CCl₂F₂ is excited at 921 cm⁻¹ in the presence of CBr₂F₂ (which does not absorb at 921 cm^{-1}), the absence of organic bromochlorides from the reaction products suggests that reaction takes place in two stages: CCl_2F_2 decomposes very quickly after excitation from a nonequilibrium molecular energy distribution, while CBr_2F_2 reacts thermally after V-T/R relaxation of the absorbed energy.

One of the hopes of infrared laser chemists is that excitation of different vibrational modes will lead to specifically different reactions. A good substrate with which to test this concept is CCl₂F₂, for which the following evidence leads to the expectation of two competing reaction channels: The laser-induced decomposition of the trifluoro analogue, CClF₃, involves a free-radical mechanism, while that of CCI₃F involves a carbene mechanism.² The corresponding primary processes for CCl_2F_2 are³

> $CCl_2F_2 \rightarrow CClF_2 + Cl_2$ $\Delta H^{\circ} = 350 \text{ kJ}$ (1)

$$\operatorname{CCl}_2F_2 \to \operatorname{CF}_2 + \operatorname{Cl}_2 \qquad \Delta H^\circ = 310 \,\mathrm{kJ}$$
 (2)

Reactions 1 and 2 are readily distinguished by the nature of the overall reaction products. In the case of (1) one would expect to find CIF₂CCCIF₂ and other free-radical products, such as $CClF_3$, while in the case of (2) one would expect to find C_2F_4 .

 CCl_2F_2 has two strong absorption bands in the CO_2 laser's tunable range.^{4,5} The band centered at 923 cm⁻¹, of b_2 symmetry, is formally an antisymmetric CCl₂ stretching mode; the band centered at 1098 cm⁻¹, of a_1 symmetry, is formally a symmetric CF₂ stretching mode. Ngai and Mann⁵ describe the dominant normal-coordinate displacements as follows: For the 923-cm⁻¹ band, C-Cl stretching, 50%; FCCl bending, 26%. For the 1098-cm⁻¹ band, C-F stretching, 59%; C-Cl stretching, 13%,

The thermal decomposition of CCl_2F_2 has been studied at 1 atm between 700 and 900 K.⁶ The products are those characteristic of a free-radical mechanism initiated by reaction 1: CClF₃, CCl₃F, and the products of their decomposition. ClF_2CCClF_2 is unstable under these conditions. There is also an unidentified white powder formed.

The ultraviolet photolysis of CCl₂F₂^{7,8} produces decomposition by both mechanisms 1 and 2. Rebbert and Ausloos⁷ find that at 213.9 nm, about 91% of the decomposition occurs via (1). However, as the wavelength of irradiation decreases to 147.0 nm, more than 80% of the decomposition occurs via (2).

Decomposition of CCl_2F_2 following excitation of the CCl_2 stretching mode with a CO_2 laser has been examined under both continuous wave^{9,10} and pulsed¹¹⁻¹³ conditions. Continuous wave irradiation results in the formation of free-radical products. Prolonged pulse irradiation, under some conditions, results in the formation of some C_2F_4 .^{11,12}

In the present paper we report on the infrared laser chemistry of CCl_2F_2 at both 921 cm⁻¹, where the CCl_2 stretching mode is excited, and at 1088 cm⁻¹, where the CF_2 stretching mode is excited. Attention will be paid to the nature of reaction products, reaction and excitation mechanism, and the amount of reaction as a function of the amount of absorbed energy. Mathematical symbols will conform to previous definitions.¹⁴

Experimental Section¹⁵

Materials. Gases were purchased from Matheson Gas Products or PCR Inc. All were of the highest purity available. All reactant gases were checked for purity by infrared and gas-liquid chromatographic (GLC) analysis, and were purified by distillations if necessary. All gases were handled and outgassed by standard vacuum-line techniques.

Reaction Cells. Most of the experiments were done in 3-cm glass cells with KCl windows. The inner cross-sectional area was 3.8 cm^2 , as compared with a beam size of $2.30-2.44 \text{ cm}^2$. Experiments using the disk calorimeter were done with special cells, with 0.014 cm^3 dead volume, described elsewhere.¹⁴

Laser. The laser system has been described in a previous paper.¹⁴ Because the laser output at 921 cm⁻¹ is relatively weak, the beam was concentrated at both frequencies from 5.8 to ~ 2.3 cm² by means of a converging lens, and then made approximately parallel by means of a diverging lens. The beam homogeneity, as measured by the fractional standard deviation σ_D , was about 13% at 1088 cm⁻¹ and 62% at 921 cm⁻¹. The high standard deviation at 921 cm⁻¹ was unavoidable because the laser beam was rectangular (long and thin rather than round). The effective pulse duration was ~ 270 ns.

Dosimetry. The accuracy of the dose measurements was improved progressively during a 2-year period. All quantitative measurements reported in this paper were made by a split-beam monitoring system, as described in Figure 2 of the previous paper.¹⁴ Meter M1 was always a pyroelectric detector. Meter M2 was either pyrolelectric or a disk calorimeter. The sensitivity of the pyrolelectric detectors was found to vary with incident energy (J) and was not uniform over the absorbing surface. When the relative positions of beam and detector were kept constant, the sensitivity V/J (where V = pyroelectric peak voltage) was represented adequately by

$$V/J = S_0/(1 + 0.178V)$$
(3)

Values of S_0 were found to be 1.63 and 1.54 V/J, respectively, for two detectors. The short-term precision of the pyroelectric measurements is about 2%. However, over longer periods we noted puzzling changes of 2-10% at the higher doses. We therefore repeated the measurements for neat CCl_2F_2 with a disk calorimeter, whose sensitivity is accurate, being determined by direct electrical calibration in each experiment.¹⁴ The two sets of measurement were quantitatively consistent.

GLC Analysis. A Varian Model 1700 instrument, equipped with a 1-mm Porapak Q column and a thermal conductivity detector, was used. Gas samples were introduced directly from the reaction cell through a Varian six-port valve.¹⁵ The detector signal was recorded and processed digitally by a Spectraphysics Autolab System I computing integrator. It was possible to resolve and determine, in the presence of largely unreacted CCl_2F_2 , the following compounds with the following accuracies: the major products CCl_3 and $sym-C_2Cl_2F_4$, 20%; the minor products CCl_3F , Cl_2FCCF_2Cl , and a cis-trans mixture of CIFC==CCIF (the two isomers are not resolved), within a factor of 2.

Infrared Analysis. A Perkin-Elmer Model 567 instrument was used. Spectra were taken directly on the reaction cells before and after laser irradiation and were recorded in different colors on a single chart, with careful attention being paid to equalizing the baselines and sensitivities. The infrared measurements were to monitor the disappearance of CCl_2F_2 , to confirm the GLC measurements for the amounts of the major products, and to confirm the absence of certain possible reaction products. For quantitative analysis, optical densities (OD) were measured at 1213 cm⁻¹ (the position of a $CClF_3$ peak), 1156 cm⁻¹ (a peak of CCl_2F_2), and 1048 and 845 cm⁻¹ (peaks of $C_2Cl_2F_4$). The actual equations were:

1213 cm ⁻¹	$L^{-1}(\text{OD}) = 0.0661y + 0.0010z$
1156 cm ⁻¹	$L^{-1}(\text{OD}) = 0.01510x + 0.0002y + 0.0002z$
$1048 + 845 \text{ cm}^{-1}$	$L^{-1} \sum (\text{OD}) = 0.0004x + 0.0004y + 0.0305z$

where L is the optical path length (in cm), $x = \Delta P_{CCl_2F_2}$, $y = P_{CCl_5}$, and $z = P_{C_2Cl_2F_4}$; pressures are in Torr. The known minor products do not absorb significantly at precisely these wavelengths.

The infrared determination of the absence of certain possible reaction products was most important in the case of C_2F_4 , whose GLC retention time is identical with that of CCIF₃. Authentic samples of C_2F_4 have a strong absorption band with peaks at 1326 and 1339 cm⁻¹. Infrared spectra taken after irradiation of neat CCl₂F₂ were completely free of any hint of absorption in this spectral range, where none of the other substances absorbs. We thus set an upper limit of 5% of the total product to the amount of C_2F_4 formed in any experiment involving neat CCl₂F₂. As a control experiment, when IR photolysis takes place in the presence of hydrogen compounds, we have no trouble at all in detecting C_2F_4 .

Results

Conversion per Flash. Results obtained for 12 Torr of CCl_2F_2 at 1088 and 921 cm⁻¹ are shown in Figure 1. As before, ¹⁴ f(%) denotes percent of conversion of reactant, per flash, calculated on the basis of the irradiated volume rather than the total volume of CCl_2F_2 . E_{abs} was measured with a disk calorimeter. The experimental points obtained at the two frequencies of irradiation fall on a single line, well within their experimental errors. Because the two frequencies were chosen so as to excite different vibrational modes, we conclude that the amount of reaction is independent of the mode of excitation, depending simply on the amount of absorbed energy per mole.

In previous work^{2,14} it was found that plots such as those of Figure 1 are approximate straight lines with negative slopes that are very close to the activation energy $E_{\rm act}$ for thermal reaction of the given substance. In the present case, the data are adequately reproduced by a straight line whose negative slope is 419 ± 40 kJ. The activation energy for thermal reaction of CCl₂F₂ has been reported⁹ as 420 kJ/mol, in good agreement, although the virtual identity of the two values is statistically fortuitous. The bond dissociation energy of the C-Cl bond, whose scission we believe to be the primary process, is 350 kJ/mol. As before,^{2,14} we conclude that reaction does *not* take place by a mechanism of laser-induced temperature jump followed by pyrolysis.

Test for Nonthermal Mechanism. To provide further proof, duplicate mixtures of 11.5 Torr of CCl_2F_2 and 5.6 Torr of SiF_4 were laser irradiated at two frequencies. At 921 cm⁻¹, where CCl_2F_2 absorbs and SiF_4 is transparent, 0.18 J/cm² was abosrbed per flash and 0.24% of CCl_2F_2 decomposed per flash. At 1031 cm⁻¹, where SiF_4 absorbs and CCl_2F_2 is transparent, 0.40 J/cm² was absorbed per flash but there was no detectable decomposition of CCl_2F_2 (<0.01% per flash). SiF_4 was inert in both experiments. If CCl_2F_2 were decomposing after the absorbed energy had become random thermal energy, the temperature jump would have been about doubled at 1031 cm⁻¹ and the decomposition of CCl_2F_2 would have been markedly greater than at 921 cm⁻¹.

Reaction Products. The two major reaction products from neat CCl_2F_2 , $C_2Cl_2F_4$ and $CClF_3$, account for $81 \pm 13\%$ of



Figure 1. Infrared laser chemistry of CCl_2F_2 at 921 cm⁻¹ (solid circles) and at 1088 cm⁻¹ (squares).



Figure 2. Total carbon of major reaction products, $C_2Cl_2F_4$ (2z) and $CClF_3$ (y) in infrared laser chemistry of CCl_2F_2 : circles, 921 cm⁻¹; squares, 1088 cm⁻¹. Data shown with "lightning bolts" were obtained by GLC analysis and pyrolelectric energy measurement and are represented approximately by the dashed line. Other data were obtained by IR analysis and disk calorimetry.

 CCl_2F_2 conversion, based on carbon balance. The identified minor products, CCl_3F and $C_2Cl_3F_3$, account at most for an additional 8%. This comprises all products we see within the



Figure 3. Ratio of major reaction products, $z = C_2Cl_2F_4$, $y = CClF_3$. Symbols for data points are the same as in Figure 2. In the upper figure, data points as a function of T_M' are shown as filled circles and filled squares.

first 10% conversion. After prolonged photolysis,¹⁵ CIFC=CCIF isomers are identified, as well as an unidentified product whose relatively long GLC retention time suggests a molecular weight greater than 150. It is clear from the above that a small (~10) percent of reaction goes to products that have not been identified. However, C_2F_4 , which one might expect to find according to a carbene mechanism (eq 2), accounts at most for half of this residue.

Within the accuracy of our data, there are no differences in either the nature or the relative amounts of the reaction products upon changing the frequency of excitation. Figures 2 and 3 show this for the two major products; z and y denote partial pressures of $C_2Cl_2F_4$ and $CClF_3$, respectively. The earlier sets of data (indicated by "lightning bolts") were obtained with pyroelectric detection, GLC analysis, and 3-cm glass cells in which the irradiated fraction φ was ~0.63. The later set of data was obtained with the disk calorimeter, infrared analysis, and a 2.6-cm Olszyna cell¹⁴ in which φ was \sim 0.76. As can be seen in Figure 2, there is a systematic difference between the two data sets. However, within each set, the results obtained at 921 cm⁻¹ mesh well with those obtained at 1088 cm⁻¹. We do not know at this time whether a systematic difference is due to the difference in detection and analysis, or whether it is a real effect caused by the difference in φ and cell length. After the rapid primary decomposition step, probably according to (1), the subsequent productforming reaction steps, at the pressures and temperatures likely to prevail, may be expected to have half-lives of at least microseconds, and probably tens of microseconds. On this time scale, the mixing of laser-heated gas with surrounding unirradiated gas must be significant, and therefore the productforming reactions probably take place at a lower effective temperature in the 3-cm cell.

6524



Figure 4. Photophysics of CCl_2F_2 in the presence of added gases: unlabeled circles, SiF4; unlabeled squares, N₂. Other data points are as follows: 1, CH₄; 2, CF₂Br₂; 3, CF₂Cl₂; 4, O₂; 5, NO.

Figure 3 shows plots of the product ratio z/y vs. E_{abs}^{-1} as well as $(RT_m)^{-1}$ and $(RT_m')^{-1}$. T_m is the temperature reached when E_{abs} adiabatically becomes random thermal energy; $T_{m'}$ is the analogous temperature when the T-jump is accompanied by perfect mixing of irradiated with unirradiated gas in the cell. It is clear that z/y increases with decreasing amount of **ab**sorbed energy, indicating that the two product-forming reactions differ in activation energy. We believe (see Discussion) that both products are formed in bimolecular reactions of CClF₂ radicals. Thus the slopes of the semilogarithmic plots vs. $(RT_m)^{-1}$ or $(RT_m')^{-1}$ indicate a probable value for the difference in activation energy of ~ 28 kJ/mol. The plot vs. E_{abs}^{-1} is included because it shows more clearly how well the product ratios obtained on excitation of the two different vibrational modes agree. The product ratios should be relatively unaffected by systematic errors.

Photophysics. Within the experimental error, E_{abs} for neat CCl_2F_2 at 12 Torr was proportional to the average dose \vec{D} in the cell, both at 921 cm⁻¹ and at 1088 cm⁻¹. E_{abs} varied from 40 to 150 kJ/mol. Average values obtained for $10^{-3}E_{abs}/\overline{D}$ (in cm² mol⁻¹) at 921 cm⁻¹ were 241 ± 12 by disk calorimetry and 244 ± 21 by pyroelectric measurement. At 1088 cm⁻¹, the corresponding average was 260 ± 26 by pyroelectric measurement. From the spectrophotometric optical densities we calculate that $10^{-3}E_{abs}/\overline{D}$ for chaotic (incoherent) radiation of low intensity is $1250 \text{ cm}^2 \text{ mol}^{-1}$ at 921 cm⁻¹ and 500 cm² mol⁻¹ at 1088 cm⁻¹. Thus we observe the interesting phenomenon that even *relative* spectrophotometric absorption coefficients may not be applied at MW/cm² coherent intensities without the risk of serious error.

Owing to mistuning of the laser, measurements with the disk calorimeter were made at 1087 cm⁻¹ rather than at 1088 cm⁻¹. $10^{-3}E_{abs}/\overline{D}$ was found to be 216 ± 16 cm² mol⁻¹, 17% smaller than at 1088 cm⁻¹. By comparison, the spectrophotometric absorbance decreases by 10%.

Effects of Added Gases. A detailed summary of our numerous experiments is given elsewhere.^{1b} We shall briefly consider some representative results, which are listed in Tables I, and II, and in Figure 4.

With the possible exception of the paramagnetic gases NO and O_2 , the addition of a second gas brings on a significant

increase in the amount of energy that is absorbed (Figure 4). However, experiments involving the addition of inert gases (N_2) and SiF₄ in Table II) show that the effectiveness at which a given amount of excitation energy can induce chemical reaction characteristically decreases. In the presence of radical trapping agents such as O2 and NO at sufficiently low pressures, the total conversion per flash actually increases, presumably due to suppression of the reverse of reaction 1, and CF₂O appears as major product (Tables I and II). The reaction of O_2 with trihalomethyl radicals, including $\cdot CCIF_2$, is a familiar reaction and results in the formation of carbohyl dihalides.^{17,18} Because the same product (CF₂O) rather than an organic nitroxide is formed in the presence of NO, we inferred from the work of Modica¹⁹ that the presence of CF₂ as primary reaction intermediate (eq 2) can be ruled out. However, owing to recently published work by Burks and Lin,²⁰ such a conclusion is now disputable and the experiment with NO is not diagnostic of reaction mechanism.

The addition of H₂ or any hydrocarbon results in the formation of C₂F₄ and/or CH₂=CF₂ as major product. That CH₂=CF₂ can be formed by a free-radical mechanism has been shown by several authors.^{21,22} C₂F₄ could result from hydrogen abstraction, CClF₂· + RH \rightarrow HCClF₂ + R, followed by thermolysis of HCClF₂²³ in the laser-heated gas.

The experiments with added HCClF₂ and CBr₂F₂ have special significance and will be taken up in the Discussion. When neat $\mathbf{CBr}_2\mathbf{F}_2$ was irradiated at 1088 cm⁻¹ where it absorbs strongly;²⁴ the only product detected by either IR or GLC analysis was BrF₂CCBrF₂. When a 12 Torr/9 Torr mixture of CCl_2F_2 and CBr_2F_2 was irradiated at 921 cm⁻¹, where CCl₂F₂ absorbs strongly and CBr₂F₂ does not, the only bromine compound detected in the product was BrF2CCBrF2. The amounts of $CBrClF_2$ and BrF_2CCClF_2 , if formed at all, were below the limits of detection (about 2% of the decrease in CCl_2F_2) by either GLC or IR. BrF₂CCClF₂ would have given a well-resolved GLC peak and practically unobscured IR absorption in the region of its very strong absorption peak at 804 cm^{-1} ²⁵ The GLC peak due to CBrClF₂ might have been obscured, but the very strong triple-peaked IR absorption band²⁴ at 866, 872, 876 cm⁻¹ should have been easily visible as a shoulder on the double-peaked CF₂Cl₂ band at 880, 887 cm^{-1} .

Discussion

Reaction Mechanism. The two most important reaction products obtained in the IR photolysis of CCl_2F_2 are $C_2Cl_2F_4$ and $CClF_3$. Equations 4-6 indicate our "best estimate" of the mechanism of their formation. The reasons for choosing these reactions rather than plausible alternatives will now be explained.

$$\operatorname{CCl}_2F_2 \rightarrow \operatorname{CCl}F_2 + \operatorname{Cl}$$
 (4) = (1)

$$2 \cdot \text{CClF}_2 \rightarrow \text{[singlet collision complex]}$$

$$\rightarrow ClF_2CCClF_2$$
 (5)

(6)

$$2 \cdot \text{CClF}_2 \rightarrow [\text{triplet collision complex}] \\ \rightarrow \text{CClF}_3 + {}^3\text{CClF}_3$$

To show that the primary step is (4) rather than (2), we argue as follows. (a) The reaction products of IR photolysis in the absence of added reactive gases are typical products for a free-radical mechanism. (b) C_2F_4 , the most likely product of a CF_2 mechanism, is not found. We believe that C_2F_4 would survive under the experimental conditions because in the IR photolysis of CCl_3F , the detected products are the three $C_2Cl_2F_2$ isomers and Cl_2 ,² showing that Cl_2 does not add to fully halogenated C=C double bonds in analogous molecules under similar conditions. (c) The major reaction product, $C_2Cl_2F_4$, which we believe to be formed by combination of

Table I. Megawatt Infrared Laser Chemistry of 12 Torr of CCl_2F_2 with Added Gases at Comparable Pressures

Added gas	Frequency, ^a cm ⁻¹	Products
None	921, 1088	$C_2Cl_2F_4 > CClF_3 > CCl_3F > C_2Cl_3F_3$
02	921	$CF_2O \gg C_2Cl_2F_4 > CClF_3$
NŌ	921	$CF_2O > C_2Cl_2F_4 > CCl_3F$; NOCl
H_2	921, 1088	$C_2F_4 \gg C_2Cl_2F_4$; HCl
CH ₄	921, 1088	$CH_2 = CF_2 \gg C_2F_4$; HCl
C_2H_4	1088	$C_2F_4 > CH_2 = CF_2 > C_2H_2$; HCl
HCCIF ₂	1088	$C_2F_4 \gg C_2Cl_2F_4$; HCl
CBr_2F_2	921	$C_2Br_2F_4 > C_2Cl_2F_4$

^a Frequency of irradiation.

•CClF₂ radicals according to (5), could conceivably be formed by insertion of CF₂ into the C-Cl bond: CF₂ + Cl-CClF₂ \rightarrow Cl•CF₂•CClF₂. To show that this reaction is unimportant, a gaseous mixture of CCl₂F₂ and HCClF₂ was irradiated at 1088 cm⁻¹, where both gases absorb similarly and strongly. The primary laser-induced reaction of the latter is HCClF₂ \rightarrow CF₂ + HCl.¹⁴ This reaction also takes place in the presence of CCl₂F₂ because we found that HCClF₂ is partly decomposed. However, the major reaction product is C₂F₄ rather than C₂Cl₂F₄ (Table I).

Concerning the mechanism of formation of CClF₃, which is the second-most important reaction product, we thought at first that CClF₃ might be formed by fluorine abstraction according to \cdot CClF₂ + CCl₂F₂ \rightarrow CClF₃ + \cdot CCl₂F. However, the following facts indicate that the analogous bromine and chlorine abstractions, which should be faster,¹⁶ are kinetically unimportant under identical or similar conditions. (a) CBrClF₂ is not formed in detectable amounts when CCl₂F₂ is photolyzed at 921 cm⁻¹ in the presence of CBr₂F₂ at comparable pressure. (b) The remarkable efficiency of C-12/C-13 isotope separation, when CCl₂F₂ is laser-irradiated at appropriate wavelengths,¹¹ could not have been achieved had there been rapid Cl-atom exchange between \cdot CClF₂ and CCl₂F₂.

We chose the mechanism shown in (6) after considering whether the formation of triplet-state products, resulting from the statistically probable triplet collisions of two free radicals, is energetically feasible. According to data in the JANAF Tables,³ standard enthalpies of formation ΔH_f° (in kJ/mol) are -182 for ¹CF₂, 12 for ³CF₂,²⁶ 239 for ¹CCl₂, and 285 for ³CCl₂.²⁶ It is clear that the triplet state of CX₂ is greatly stabilized, relative to the singlet state, by substitution of Cl for F. On adopting a mean value of those for ³CF₂ and ³CCl₂, we obtain $\Delta H_f^{\circ} = 148$ kJ/mol for ³CClF, and thus arrive at the following thermochemistry:

 $2 \cdot \text{CClF}_2 \xrightarrow{\mathsf{T}} \text{CClF}_3 + {}^3\text{CClF} \quad \Delta H^\circ = -23 \text{ kJ} \quad (6a)$ $2 \cdot \text{CClF}_2 \xrightarrow{\mathsf{T}} \text{CCl}_2 F_2 + {}^3\text{CF}_2 \qquad \Delta H^\circ = +50 \text{ kJ} \quad (7)$

It appears that for triplet radical-radical reaction, C-F abstraction is exothermic while C-Cl abstraction is endothermic. If the stated mechanism (4)-(6) is correct, the difference in activation energy, $E_6 - E_5$, is estimated from the upper plots in Figure 3 as ~28 kJ.

We are currently planning experiments to test the presence of ³CCIF and examine its eventual conversion to stable products more decisively. According to the mechanism stated in eq 4-6, the disappearance of $CCl_2F_2(-x)$ should be related to the formation of $C_2Cl_2F_4(z)$ and $CCIF_3(y)$ by -x = 2y + 2z. The experimental ratio of -2(y + z)/x is 0.96 \pm 0.12, essentially unity.

Excitation Mechanism. It is clear, from the good correspondence of results obtained at 921 and at 1088 cm^{-1} , that

 Table II. Semiquantitative Photochemistry and Photophysics at 921 cm⁻¹ for CF₂Cl₂ (12 Torr)

Added gas, Torr	$\frac{10^{-3} E_{abs}/\overline{D}.^{a}}{cm^2 mol^{-1}}$	$E_{abs} = 90 \text{ kJ/mol}^{b}$ $\frac{f}{f_0} \frac{z}{z_0}$	
None	242	1.00	1.00
$N_2(22.1)$	317	0.1	0.03
$O_2(5.2)$	248	2.2	0.16
(7.9)	266	0.8	0.04
NO(8.4)	232	2.2	0.73
$SiF_4(4.8)$	261	0.25	0.25
CH ₄ (9.4)	279	-	-
$\operatorname{CBr}_2\operatorname{F}_2(9.2)$	286	0.6°	0.83 °

^{*a*} Average values for 60 < E_{abs} < 120 kJ/mol; typical standard deviation is ±20. ^{*b*} Obtained by interpolation. E_{abs} measured with pyroelectric detectors; f/f_0 and z/z_0 measured by GLC; $z = C_2Cl_2F_4/flash$. ^{*c*} (quasi- E_{abs})/ $E_{abs} = 0.94_5$ (from f/f_0), 0.98₅ (from z/z_0); average 0.97.

the laser-activated molecules at the moment of reaction retain practically no memory of the original mode of excitation. The sequence of microscopic events leading up to primary reaction is of course of interest. We shall describe an experiment whose outcome is particularly revealing.

When 12 Torr of CCl_2F_2 is irradiated at 921 cm⁻¹ in the presence of 9 Torr of CBr_2F_2 , energy is absorbed only by CCl₂F₂, yet both substances react to yield free-radical products. CCl₂F₂ gives C₂Cl₂F₄ as major product, CBr₂F₂ gives $C_2Br_2F_4$ as the only detected product, suggesting the intermediate formation of both \cdot CClF₂ and \cdot CBrF₂. Yet significantly, there is no trace of the bromochloride ClF_2CCBrF_2 , whose formation would have been surely expected had the two free-radical species been present simultaneously. Thus the reactions of CCl_2F_2 and CBr_2F_2 appear to be time resolved. We believe that CCl_2F_2 (which absorbs at 921 cm⁻¹) decomposes first, while CBr₂F₂ decomposes later, after an appropriate lag period to allow for intermolecular energy transfer, and mostly after the temperature jump. Because of the known high rate constants for the combination of halomethyl radicals,¹⁶ it is plausible to propose that the CClF₂ radicals will have largely decayed to stable products during the lag period, before the formation of CBrF2 radicals is well underway. After the temperature jump, one would expect the thermal reaction of CBr₂F₂ to be faster by several orders of magnitude than that of CCl_2F_2 , because of the lower bond dissociation energy, 265 kJ/mol for C-Br vs. 350 kJ/mol for C-Cl. As far as we know, this is the first instance of *chemical* proof that the reaction of laser-activated CCl_2F_2 molecules is photochemical rather than thermal.

Further insight is provided by the amount of reaction of CCl_2F_2 . As shown by typical results in Table II, when $E_{abs} = 90 \text{ kJ/mol}$, conversion per flash (f) and amount of $C_2Cl_2F_4$ produced per flash (z) in the presence of CBr_2F_2 are only slightly smaller than the corresponding values, f_0 and z_0 , obtained when $E_{abs} = 90 \text{ kJ/mol}$ in the absence of CBr_2F_2 . Because the magnitude of the T-jump of the laser-heated gas will be considerably smaller in the presence of CBr_2F_2 , the near absence of any effect due to CBr_2F_2 on f and z further confirms the previous conclusion that the primary step for CCl_2F_2 reaction takes place before the T-jump is well under way. Expressed in terms of E_{abs} , at 90 kJ/mol in the presence of CBr_2F_2 , laser-activated CCl_2F_2 decomposes essentially with the same yield and chemistry per flash as it would in the absence of CBr_2F_2 at $87 \pm 2 \text{ kJ/mol}$ (Table II, footnote c).

This result becomes revealing of microscopic events when analyzed in terms of number of gas-kinetic collisions. In a gas phase consisting of 12 Torr of CCl_2F_2 and 9 Torr of CBr_2F_2 , the average number of collisions experienced by a CCl_2F_2 molecule during the 270 ns effective duration of the laser pulse

is ~80 with CCI_2F_2 molecules and ~60 with CBr_2F_2 molecules. The collision number, Z_{10} , for V-T/R relaxation from the lowest excited vibrational state to the ground state of CCl_2F_2 is 212 collisions,²⁷ so that the laser pulse duration is about one half-life. By comparison, the maximum percent of V-T/R energy transfer for CCl_2F_2 in the actual experiment is $3 \pm 2\%$.²⁸ As is well known (and expressed mathematically by the Lambert-Salter rule²⁹), collisional V-T/R energy transfer is efficient only from states of low vibrational excitation energy. Our results therefore indicate that, prior to reaction, at most a small fraction of the vibrational excitation energy resides in low-energy modes.

Even more remarkable is the evident unimportance of V-V'energy transfer between CCl_2F_2 and CBr_2F_2 . These two molecules are as nearly alike (in structure, symmetry, and chemical type) as two distinct molecules can be, short of being isotopic variants of the same molecular species. The normalmode frequencies for analogous vibrations range from fairly similar to near resonant.24 In particular, the two C-F stretching modes of CCl₂F₂ at 1101 and 1159 cm⁻¹ are within 20 cm^{-1} of the analogous normal-mode frequencies of CBr_2F_2 . The C-Cl stretching mode at 667 cm^{-1} is within 50 cm^{-1} of the analogous C-Br stretching mode. The torsional mode at 318 cm⁻¹ is within 20 cm⁻¹ of a bending mode of CBr₂F₂. If the excitation energy of CCl₂F₂ were residing in, or exchanging rapidly into, any of these modes, one might reasonably expect³⁰ that more than a maximum of $3 \pm 2\%$ of the energy is transferred to vibrational modes of CBr₂F₂. The fact that it does not happen indicates that the vibrational energy of CCl_2F_2 , prior to reaction, remains narrowly localized.

Following earlier analysis,^{2,3} we conclude from the agreement of the slope of Figure 1 with the experimental activation energy for thermolysis, that the laser-induced reaction takes place from a molecular energy distribution which closely resembles a Boltzmann distribution of energy in a single vibrational mode. Let us suppose that resemblance mirrors reality. Then what is the frequency of the single mode? In the present experiment the excitation frequency is 921 cm^{-1} . We have seen that the energy is not likely to reside in the two higher-frequency modes of CCl₂F₂, because of the unimportance of energy transfer to CBr₂F₂. Nor is it likely to reside in several of the lower frequencies, for the same reason, and also because of the unimportance of V-T/R energy transfer. This does not leave much choice, and the simplest choice is that the energy remains localized in the original mode of excitation, perhaps because the molecules are "locked" coherently onto the laser frequency.2,31

Finally, it must be pointed out that if this explanation can be accepted, the molecules, while locked onto the laser frequency, undergo vibrational motions which are different from the kind of motion normally called translation along the reaction coordinate. However, at some time shortly after the laser pulse, the frequency lock becomes sufficiently weak that the molecules become able to escape from the forced mode of motion and, at constant energy, explore other regions of phase

space. Those molecules which possess enough energy to react will, in due course, find a reaction channel and go on to products. Judging by the Monte Carlo calculations of Bunker and Hase³² for a somewhat more complex process, the isomerization of CH₃NC, the time required for finding a reaction channel is of the order of 10^{-9} to 10^{-11} s, starting at a random point in phase space, at energies comparable to those required for reaction here. By comparison, the mean time between successive (gas-kinetic) collisions of a CCl_2F_2 molecule, under the present conditions, is $\sim 2 \times 10^{-9}$ s. Thus, at the end of the period of forced motion, most molecules with enough energy to react will react before the next collision occurs, regardless of their mode of excitation. The model can therefore explain the observation, in the case of CCl_2F_2 , that the amount of reaction depends only on E_{abs} and is independent of the mode of excitation.

References and Notes

- (1) (a) Work supported by the National Science Foundation. (b) Abstracted from the Ph.D. dissertation of Gregory A. Hill, Brandeis University, 1977. Present address: Department of Chemistry, Columbia University, New York, N.Y. 10027.
- (2) D. F. Dever and E. Grunwald, J. Am. Chem. Soc., 98, 5055 (1976). (3) Thermochemical data are based on the JANAF Tables, Natl. Stand. Ref. Data. Ser., Natl. Bur. Stand., No. 37 (1971),
- (4) H. W. Thompson and R. B. Temple, J. Chem. Soc., 1422 (1948).
- (5) L. H. Ngai and R. H. Mann, J. Mol. Spectrosc., 38, 322 (1971)
- (6) A. B. Trenwith and R. H. Watson, J. Chem. Soc., 2369 (1957)
- R. E. Rebbert and P. J. Ausloos, J. Photochem., 4, 419 (1975)
 R. Milstein and F. S. Rowland, J. Phys. Chem., 79, 669 (1975)
- (9) M. P. Freeman, D. N. Travis, and M. F. Goodman, J. Chem. Phys., 60, 231 (1974)
- (10) R. N. Zitter, R. A. Lau, and K. S. Wills, J. Am. Chem. Soc., 97, 2578 (1975)
- (11) J. J. Ritter and S. M. Freund, J. Chem. Soc., Chem. Commun., 811 (1976)
- (12) J. L. Lyman and S. D. Rockwood, J. Appl. Phys., 47, 595 (1976).
- (13) R. V. Ambartzumian, N. V. Chekalin, V. S. Letokhov, and E. A. Ryabov, Chem. Phys. Lett., 36, 301 (1975).
- (14) E. Grunwald, K. J. Olszyna, D. F. Dever, and B. Knishkowy, J. Am. Chem. Soc., preceding paper in this issue. (15) For full experimental details, see ref 1b.
- (16) For examples, see D. C. Nonhebel and J. C. Walton, ''Free-Radical Chemistry", Cambridge University Press, Cambridge, England, 1974.
- (17) D. Marsh and J. Heicklen, J. Phys. Chem., 69, 4410 (1965).
- R. Milstein and F. S. Rowland, J. Phys. Chem., 79, 669 (1975).
 A. P. Modica, J. Chem. Phys., 46, 3663 (1967).
- T. L. Burks and M. C. Lin, J. Chem. Phys., 64, 4235 (1976) (20)
- (21) R. D. Giles and E. Whittle, Trans. Faraday Soc., 61, 1425 (1965).
- 22) J. M. Farrar and Y. T. Lee, Annu. Rev. Phys. Chem., 25, 376 (1974). (23) G. R. Barnes, R. A. Cox, and R. F. Simmons, J. Chem. Soc. B, 1176 (1971)
- (24) E. K. Plyler and N. Acquista, J. Res. Natl. Bur. Stand., 48, 92 (1952); E. K. Plyter and W. S. Benedict, ibid., 47, 212 (1951).
- (25) R. E. Kagarise, J. Chem. Phys., 26, 380 (1957).
- (26) ΔH₄^o for ¹CF₂ as listed in ref 3, plus the result of a careful quantum-me-chanical calculation (which agrees with an estimate in ref 3) for the trip-let-singlet energy difference. V. Staemmler, *Theor. Chim. Acta*, **35**, 309 (1974). The value given for ³CCl₂ is based on an estimate in ref 3.
- (27) T. D. Rossing and Š. Legvold, *J. Chem. Phys.*, 23, 1118 (1955).
 (28) We are comparing the value of *quasi-E_{abs}* of 87 ± 2 kJ/mol, based on the amount of chemistry per flash, with the actual *E_{abs}* of 90 kJ/mol.
- (29) J. D. Lambert and R. Salter, Proc. R. Soc. London, Ser. A, 253, 277 (1959).
- (30) E. Weitz and G. Flynn, Annu. Rev. Phys. Chem., 25, 303 (1974).
 (31) E. Grunwald and K. J. Olszyna, Laser Focus, 12 (6), 41 (1976).
- (32) D. L. Bunker and W. L. Hase, J. Chem. Phys., 59, 4621 (1973).